

AN ANOMALOUS COMPONENT OF GALACTIC EMISSION

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ABSTRACT

We present results from microwave background observations at the Owens Valley Radio Observatory. These observations, at 14.5 and 32 GHz, are designed to detect intrinsic anisotropy on scales of 7'–22'. After point-source removal, we detect significant emission with temperature spectral index $\beta \simeq -2$ toward the north celestial pole (NCP). Comparison of our data with the *IRAS* 100 μm map of the same fields reveals a strong correlation between this emission and the infrared dust emission. From the lack of detectable H α emission, we conclude that the signals are consistent either with flat-spectrum synchrotron radiation or with free-free emission from $T_e \gtrsim 10^6$ K gas, probably associated with a large H I feature known as the NCP Loop. Assuming $\beta = -2.2$, our data indicate a conversion $T_f/I_{100\mu\text{m}} = 7.5 \times 10^{-2} \nu_{\text{GHz}}^{-2.2} \text{ K (MJy sr}^{-1})^{-1}$.

The detection of such a component suggests that we should be cautious in any assumptions made regarding foregrounds when designing experiments to map the microwave background radiation.

Subject headings: cosmic microwave background — dust, extinction — H II regions — supernova remnants

1. OBSERVATIONS

Since 1993, the Owens Valley Radio Observatory (OVRO) 5.5 m telescope (see Herbig et al. 1995) has been used for extensive observations at 32 GHz of the cosmic microwave background (CMB) on angular scales of 7'–22'. The receiver input is switched at 500 Hz between two beams of 7:35 (FWHM) separated by 22:16. The OVRO 40 m telescope, underilluminated at 14.5 GHz to match the 5.5 m beam, provides a second frequency channel for spectral discrimination of foregrounds (see Table 1). Both receivers detect right-circular polarization.

In the “RING5M” experiment, we observe 36 fields at $\delta \simeq 88^\circ$ spaced by the 22:16 beam throw in a ring around the north celestial pole (NCP) (see Fig. 1). To minimize variations in differential contributions from the ground, each field is observed only within ± 20 minutes of upper culmination. In each flux measurement, the telescope moves in azimuth to alternate the beams on the main field; each measurement is thus the difference between the signal from the main field and the average of the signals in the two adjacent fields (Readhead et al. 1989). As a result, a strong signal in one field produces a negative signal half as strong in each of the two flanking fields (see Fig. 2).

To estimate the contribution of point sources, the RING5M fields were mapped on the VLA at 8 GHz to a sensitivity of 0.25 mJy; a total of 34 sources were detected with $S_{8\text{GHz}} \gtrsim 2$ mJy. Subsequent monthly VLA monitoring of these sources at 8 and 15 GHz provided accurate measurements of flux densities and spectral indices, enabling us to estimate the flux densities at both 14.5 GHz and 32 GHz and to subtract the point-source contribution from these data sets.

We report here only on results relevant to an anomalous foreground we have detected. The implications for CMB observations and a full discussion and analysis of our results will be presented in subsequent papers.

2. ANALYSIS

Our cumulative observations of the RING5M fields have achieved a 1σ rms sensitivity of 17 μK per field at 32 GHz, and 15 μK per field at 14.5 GHz. Signals $\geq 200 \mu\text{K}$ are seen in both channels, and excellent reproducibility of these data between the 1994, 1995, and 1996 observing seasons indicates that they represent real structure on the sky.

In addition to good agreement between independent data sets at the same frequency, a Spearman rank test (Kendall 1955; Press et al. 1988), modified to account for correlations introduced by the switching (see § 3), finds a correlation $r_s = 0.84$ between the two frequencies, with a significance $P(r_s \geq 0.84) = 7 \times 10^{-7}$. The strength of the observed correlation between independent observations on separate telescopes is further evidence that the signals are astronomical in origin and not artifacts of the observing procedure.

Since the only common element between the two channels is our observing strategy, we explored the possibility of systematic contamination by observing the RING5M fields at 14.5 GHz for two weeks at *lower* culmination. The data obtained showed the same signals, to within the observed 80 μK scatter between 2 week subsets of our upper culmination data.

In all of the following discussion, we use data sets from which the point-source contributions have been subtracted. Apart from one source, which affected three of our 36 fields, the contributions of point sources were much smaller than the detected signals. The maximum 1σ error on the estimated point-source contribution to a field was 22 μK . There is therefore no doubt that we have detected a significant astronomical signal that is not due to point sources.

2.1. Spectral Index of the Foreground

The strongest signals seen in both the 14.5 and 32 GHz channels have steep spectral indices and amplitudes $T \sim 300 \mu\text{K}$ at 14.5 GHz. A likelihood analysis yields $\beta = -1.1^{+0.2}_{-0.3}$ for the spectral index of the data set as a whole, indicating the presence of significant emission with $\beta < 0$.

Recent Westerbork observations (Wieringa et al. 1993) reveal polarized structures at high Galactic latitude as bright

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TABLE 1
PARAMETERS FOR THE OVRO 5.5 AND 40 METER TELESCOPES

Parameter	5.5 m	40 m
Frequency.....	32 GHz	14.5 GHz
Bandwidth.....	5.6 GHz	3 GHz
System temperature.....	30 K	48 K
Beam efficiency.....	0.61	0.70
Beamwidth (FWHM).....	7:35	7:35
Beam throw.....	22:16	21:50
Main-beam solid angle.....	5.12×10^{-6} sr	5.38×10^{-6} sr

as 8 K at 325 MHz. These features, on scales of $4'$ – $30'$, are seen in linear polarization only; the corresponding total intensity maps are extremely smooth, and upper limits less than 0.5–1 K are set on any counterparts in total intensity. This is interesting, as approximately 8 K features with spectral index $\beta = -2.7$ can just reproduce the observed structures at 14.5 GHz if we were 100% sensitive to linear polarization. Tests of our 14.5 GHz polarizers, however, indicate less than 6% contamination from linear polarization across our bandpass, so it is clear that such polarized features cannot account for the signals we have detected.

Moreover, given the smoothness of the total intensity maps, it is highly improbable that the structure in the polarized emission is due to variations in intrinsic polarization angle, and the polarized structure is interpreted as Faraday rotation of an intrinsically smooth, polarized synchrotron background by an intervening screen. As a result, the polarization angle will have the ν^{-2} dependence of Faraday rotation and the structure should be negligible at 14.5 GHz.

Total intensity maps from the WENSS survey (de Bruyn

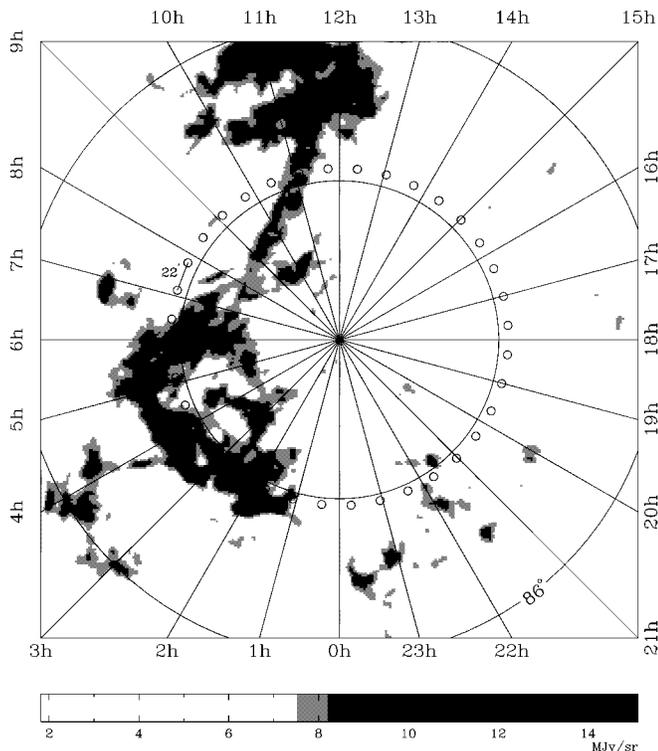


FIG. 1.—*IRAS* 100 μm map in J2000 coordinates, with RING5M fields overlaid, ranging from $b \approx 25^\circ$ to $b \approx 30^\circ$ in Galactic latitude. The spacing of the fields is about 22', and the FWHM of the beam is 7:35.

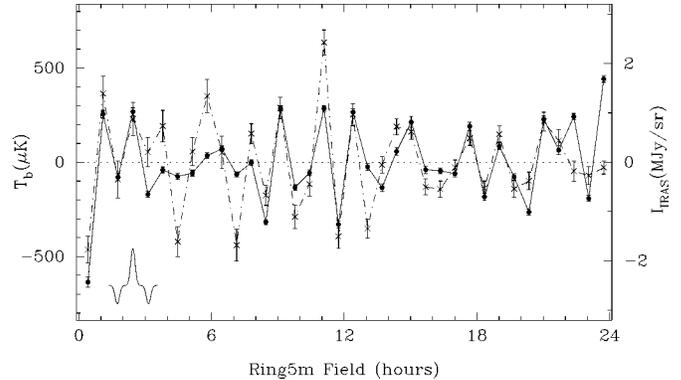


FIG. 2.—Comparison of the 14.5 GHz data (solid line) in μK , with the *IRAS* 100 μm convolution (dot-dashed line). Errors for the *IRAS* data points are the estimated standard deviation of the convolution. The dotted line essentially coincident with the x -axis is the free-free signal, in μK , inferred from H α images of the NCP fields, assuming $T_e = 10^4$ K. At bottom left is the “triple beam” pattern due to the double switching.

1996), covering 20 of the 36 RING5M fields, show no detectable signals after removal of discrete sources. Comparison of the maximum signal at 14.5 GHz with the rms from the WENSS maps in the overlap fields places a lower limit $\beta \geq -2.2$ on the spectral index of the foreground we have detected.

Thus, based on the WENSS maps, we know that the contribution of any steep-spectrum ($\beta < -2.2$) component is negligible, and we now investigate what foreground spectral index is favored by our data. We model our data as a Gaussian CMB component in the presence of a single foreground of variable strength but constant spectral index β . Defining $\Delta T_i = \delta T_i - \frac{1}{2}(\delta T_{i+1} + \delta T_{i-1})$, as described in § 1, for each field i , we measure

$$\Delta T_i(\nu) = \Delta T_{\text{cmb}} + \Delta T_{\text{fore}} \nu^\beta. \quad (1)$$

Given $\Delta T_i(\nu)$ measured at two frequencies ν_1 and ν_2 , we can solve for the CMB component in terms of the unknown spectral index of the foreground:

$$\Delta T_{\text{cmb}}(\beta) = \frac{\Delta T_i(\nu_1) \nu_1^{-\beta} - \Delta T_i(\nu_2) \nu_2^{-\beta}}{\nu_1^{-\beta} - \nu_2^{-\beta}}. \quad (2)$$

The likelihood function for the CMB component (see, for example, Readhead et al. 1989) is then given by

$$L(\sigma_{\text{cmb}}, \beta) = \prod_{i=1}^{36} \frac{1}{\sqrt{2\pi(\epsilon_i^2 + \sigma_{\text{cmb}}^2)}} \exp \left[-\frac{\Delta T_{\text{cmb}}^2(\beta)}{2(\epsilon_i^2 + \sigma_{\text{cmb}}^2)} \right], \quad (3)$$

where ϵ_i is the error in the residual CMB component, and σ_{cmb}^2 is the intrinsic CMB variance. The likelihood constructed from the point-source subtracted data sets at 14.5 and 32 GHz peaks at $\beta = -2.25$, with $\beta > -1.8$ ruled out at the 1σ level. Our data are thus consistent with a foreground of spectral index $\beta \sim -2$, and we conclude that the foreground is either unusually flat-spectrum synchrotron radiation or free-free emission.

3. *IRAS* 100 MILLIMETER MAPS

In an attempt to correlate this component with known Galactic foregrounds, we convolved the *IRAS* 100 μm map (*IRAS* Explanatory Supplement 1988) of the NCP with our beam and beam switch. The source-subtracted 14.5 GHz data,

along with results of the *IRAS* convolution, are shown in Figure 2.

We find a clear correlation between the *IRAS* 100 μm maps and our 14.5 GHz data set. To assess the significance of this result without a priori knowledge of the distribution of the *IRAS* 100 μm brightness or 14.5 GHz temperature on 7' scales, we use Spearman's rank correlation coefficient r_s . Since this depends only on the data *ranks*, whose distribution is known, and not on the values themselves, we can determine the significance of an observed value of r_s unambiguously. The observed correlation between the 14.5 GHz and *IRAS* 100 μm data is $r_s = 0.73$, and, for 36 independent fields, the probability of observing $r_s \geq 0.73$ by chance is $P(r_s \geq 0.73) = 4.5 \times 10^{-7}$. Because of our switching strategy, however, only every third field is actually independent, and numerical simulations show that this reduces the significance to 6.7×10^{-5} .

We note that the region spanning $3^{\text{h}}-8^{\text{h}}$, where the correlation is weakest, is also the region where we see the strongest signals at 32 GHz; the spectral indices of these fields are consistent with $\beta = 0$, indicating the presence of a significant CMB signal.

Taking $\beta = -2.2$ as the spectral index of the foreground most consistent with both our data and the WENSS maps, we can use the 14.5 GHz and 32 GHz data to solve for the foreground component in the manner of equation (2). A linear fit to the foreground component yields the conversion

$$T_f/I_{100\mu\text{m}} = 7.5 \times 10^{-2} \nu_{\text{GHz}}^{-2.2} \text{ K (MJy sr}^{-1}\text{)}^{-1}. \quad (4)$$

4. H α OBSERVATIONS OF THE NORTH CELESTIAL POLE

Gaustad, McCullough, & Van Buren (1996) have recently estimated the free-free contamination of small-scale anisotropy experiments through H α observations of the NCP. For the brightness temperature of optically thin hydrogen, we expect

$$T_b(\mu\text{K}) = \frac{5.43}{\nu_{10}^2 T_4^{1/2}} g_{\text{ff}} \text{EM}_{\text{cm}^{-6}\text{pc}}, \quad (5)$$

where

$$g_{\text{ff}} = 4.69(1 + 0.176 \ln T_4 - 0.118 \ln \nu_{10}) \quad (6)$$

is the free-free Gaunt factor (Spitzer 1978), the frequency is $10\nu_{10}$ GHz, the electron temperature $T_e = 10^4 T_4$ K, and $\text{EM} \equiv \int n_e^2 dl$ is the emission measure. For $T_e \leq 2.6 \times 10^4$ K, the H α surface brightness in rayleighs ($1 \text{ R} = 2.42 \times 10^{-7} \text{ ergs cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$ at H α) is given by

$$I_{\text{H}\alpha}(\text{R}) = 0.36 \text{EM}_{\text{cm}^{-6}\text{pc}} T_4^{-0.9} \quad (7)$$

(Kulkarni & Heiles 1987).

Simulating our observing procedure on the maps of Gaustad et al. (1996), we measure $\langle \Delta I \rangle_{\text{rms}} \leq 0.1 \text{ R}$ on 7' scales in H α . If we assume $T_e = 10^4$ K for the temperature of the emitting gas, then the inferred upper limit on the rms at 14.5 GHz due to free-free emission is $\langle \Delta T_{\text{ff}} \rangle_{\text{rms}} \leq 3.2 \mu\text{K}$, a factor of about 60 lower than the observed $\langle \Delta T_{\text{ff}} \rangle_{\text{rms}} = 203 \mu\text{K}$. Furthermore, the H α maps are featureless; in the 36 RING5M fields, no signals are seen with $|\Delta I| > 0.2 \text{ R}$. For the approximately 300 μK signals we detect, equations (5), (6), and (7) predict an H α brightness $|\Delta I| \sim 9 \text{ R}$.

If considerable dust lies along the line of sight to the NCP, extinction might account for the low levels of observed H α emission; estimates from the *IRAS* 100 μm intensities, how-

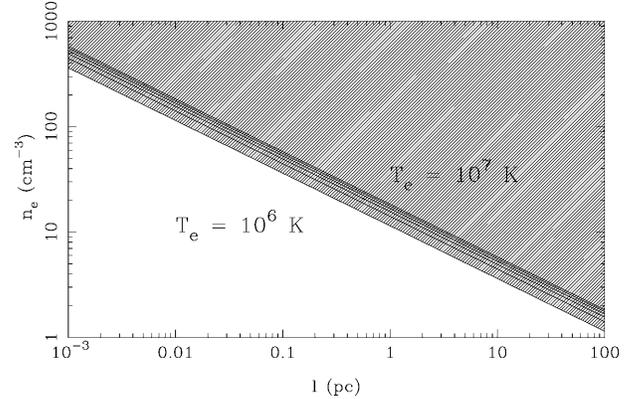


FIG. 3.—Allowed $n_e - l$ parameter space (shaded region) for $T_e \geq 10^6$ K. We can exclude $l > 100$ pc, as this requires that inhomogeneities be aligned along the line of sight to within about 6×10^{-3} rad, since we see fluctuations on the scale of the 22' beam throw. The solid lines correspond to $T_e = 10^6, 2 \times 10^6, 4 \times 10^6, 6 \times 10^6, 8 \times 10^6,$ and 10^7 K.

ever, imply $\lesssim 0.6$ mag of visual extinction (Simonetti, Dennison, & Topansa 1996), so that the upper limits on free-free emission can be increased by 74% at most.

As T_e is increased beyond 2.6×10^4 K, the allowed orbital space for recombination shrinks, and equation (7) is no longer valid; for $T_e > 2.6 \times 10^4$ K, a fit to the H α recombination coefficient gives $\alpha_{\text{H}\alpha} \propto T^{-1.2}$ (Ferland 1980). The presence of 300 μK free-free emission can therefore be reconciled with the observed 3σ H α limit if the emission is due to gas at $T_e \gtrsim 10^6$ K. For these temperatures, free-free brightness temperatures of 300 μK at 14.5 GHz require an EM $\gtrsim 131$. The corresponding allowed $n_e - l$ parameter space is shown in Figure 3.

5. DISCUSSION

The observed structure in the *IRAS* 100 μm map of the NCP region is part of a large H I feature known as the NCP Loop. This feature, which encompasses all of the 36 RING5M fields, has been modeled by Meyerdieks, Heithausen, & Reif (1991) as the wall of an expanding cylindrical shock. While the production of a dense ionized component such as that implied by Figure 3 may pose significant difficulties—such structures will be extremely overpressured and must necessarily be transient phenomena—it is intriguing that the combination of large emission measure and high temperature arrived at by interpreting the observed structure at 14.5 GHz as free-free emission is suggestive of just such a shocked component of the interstellar medium (ISM).

If the emission is due to $\gtrsim 10^6$ K gas, then this component should have a counterpart in soft X-rays; the absence of any *ROSAT* PSPC pointings near the NCP and the low resolution of available all-sky surveys, however, prevent any useful comparison with existing data sets.

Perhaps more plausible, though no less anomalous, is the possibility that the observed structure at 14.5 GHz is due to flat-spectrum synchrotron emission. Synchrotron spectral indices as flat as $\beta = -2.0$ are typically observed only in plerions associated with the very youngest supernova remnants (SNRs) (Green 1996), and we would not expect such emission from the NCP Loop—an old remnant, with expansion velocity $v \simeq 20 \text{ km s}^{-1}$. A notable exception to the general steepness of Galactic synchrotron radiation, however, is the filamentary structure observed toward the Galactic center. These features, consisting of long, nearly one-dimensional threads, have spec-

tral indices $-2.2 \leq \beta \leq -1.9$, yet they show considerable linear polarization, suggesting that the dominant emission mechanism is synchrotron (Yusef-Zadeh 1989). Although such structures would suffer from a similar lifetime problem as free-free filaments and require recent injection of high-energy electrons to maintain a flat spectrum, they would obviate the high temperature and pressure required in the case of free-free emission.

6. CONCLUSIONS

We have detected a significant correlation between emission with temperature spectral index $\beta \sim -2$ observed at 14.5 GHz in the RING5M experiment and the *IRAS* 100 μm maps. If this is free-free emission, then the lack of accompanying $\text{H}\alpha$ emission implies that it is from a component of the ISM with $T_e \gtrsim 10^6$ K. The large EM required to produce the observed signals at these temperatures is typical of SNR, an interpretation supported by the morphology of the NCP Loop, with which the *IRAS* emission is associated.

Kogut et al. (1996) have recently reported a large angular scale correlation of the residual free-free component in the *COBE* Differential Microwave Radiometer (DMR) sky maps with far-infrared DIRBE emission. The level of this signal at 53 GHz, however, is consistent with predictions from $\text{H}\alpha$ observations, implying that on 7° scales, the observed free-free emission is from a $T_e \sim 10^4$ K phase of the ISM. Moreover, if

the correlation with free-free emission persists to small scales, then the power spectrum of the high-latitude DIRBE 240 μm maps $P(l) \propto l^{-3}$, where $l \sim 60/\theta^\circ$, implies a level of free-free at 0.1° scales marginally consistent with the limit inferred from $\text{H}\alpha$ observations.

If the observed foreground is not unique to the NCP region, then our results imply that such emission could be a serious contaminant to small-scale CMB measurements in other areas of sky. This component does, however, have a significantly steeper spectral index and may be subtracted out by multifrequency observations. Moreover, further observations now in progress to determine the extent of the correlation between 14.5 GHz and 100 μm emission indicate that these results for the NCP region are atypical.

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